

Kinetic Theory of the Dyakonov-Shur Instability

joint work with Andrew Lucas, Stanford University

arXiv:1801.01501

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Spontaneous generation of terahertz radiation in a 2DEG

Original proposal by Dyakonov and Shur (1993) (implicitly $\tau_{ee} \ll 1$):

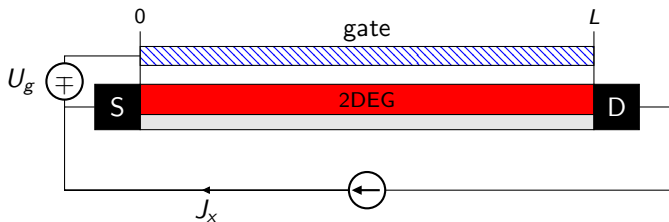


Figure: Schematic experimental setup

Idea: spontaneous fluctuations “blow up” in time: $\propto e^{\gamma t}$ with $\gamma > 0$

Non-standard boundary conditions:

$$\delta n(x=0) = 0, \quad \delta J(x=L) = 0$$

This study: transition **hydrodynamic** ($\tau_{ee} \ll 1$) to **ballistic** ($\tau_{ee} \gg 1$)

Boltzmann kinetic equation

2D Fermi liquid, fermion distribution function $f(\mathbf{x}, \mathbf{p}, t)$

$$\partial_t f + \frac{\mathbf{p}}{m} \cdot \nabla_{\mathbf{x}} f + \mathbf{F}^{\text{ext}} \cdot \nabla_{\mathbf{p}} f = \mathcal{C}[f]$$

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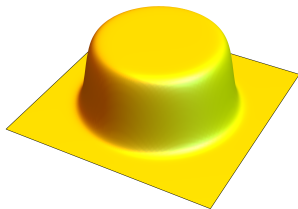
BGK approximation for collision operator:

$$\mathcal{C}[f] = -\frac{f - F_{\text{eq}}[f]}{\tau_{\text{ee}}}$$

with

$$F_{\text{eq}}[f] \equiv \Theta(\mu[f] + \mathbf{u}[f] \cdot \mathbf{p} - \epsilon(\mathbf{p})),$$

$\mu[f]$ and $\mathbf{u}[f]$ determined self-consistently
by density and momentum conservation



Guo et al. 2016, 2017; Jong and Molenkamp 1995; Lucas 2017

Boltzmann kinetic equation

2D Fermi liquid, fermion distribution function $f(\mathbf{x}, \mathbf{p}, t)$

$$\partial_t f + \frac{\mathbf{p}}{m} \cdot \nabla_{\mathbf{x}} f + \mathbf{F}^{\text{ext}} \cdot \nabla_{\mathbf{p}} f = C[f]$$

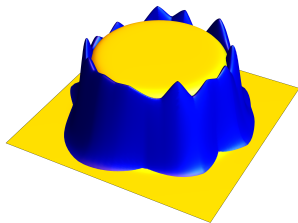
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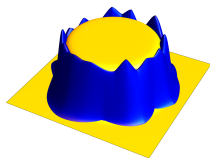
Guo et al. 2016, 2017; Jong and Molenkamp 1995; Lucas 2017

Boltzmann equation: harmonic representation

$$f(\mathbf{x}, \mathbf{p}, t) \approx F_{\text{eq}}(\mathbf{p}) + \delta(\epsilon(\mathbf{p}) - \mu)\Phi(\mathbf{x}, \mathbf{p}, t)$$

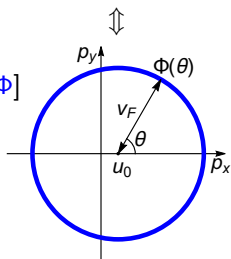
Harmonic representation

$$\Phi(\mathbf{x}, \mathbf{p}, t) = e^{-i\omega t} \sum_{n \in \mathbb{Z}} a_n(\mathbf{x}) e^{in\theta}$$



Inserted into Boltzmann equation \rightsquigarrow

$$i\omega\Phi = (u_0 + v_F \cos(\theta))\partial_x\Phi + \frac{2v_g^2}{v_F} \cos(\theta)\partial_x a_0 + \frac{1}{\tau_{ee}}P[\Phi]$$



$$u_0 = J_x / \rho \quad \left| \begin{array}{l} \text{fluid velocity} \\ v_F \\ \text{Fermi velocity} \end{array} \right.$$

$$v_g = \sqrt{\frac{e^2 n_0}{mC}} \quad \left| \begin{array}{l} \text{gate "velocity"} \end{array} \right.$$

$$P[\Phi] = \sum_{|n| \geq 2} a_n e^{in\theta} \quad \left| \begin{array}{l} \text{collision operator (projection)} \end{array} \right.$$

Boundary conditions

Dyakonov-Shur boundary conditions:

$$\delta n(0) = 0 \quad \Leftrightarrow \quad a_0(0) = 0,$$

$$\delta J(L) = 0 \quad \Leftrightarrow \quad u_0 a_0(L) + \frac{v_F}{2}(a_1(L) + a_{-1}(L)) = 0$$

Additionally:

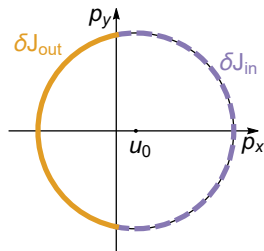
“Clean” (specular reflection)
boundary conditions:

$$\Phi(\theta) = \Phi(\pi - \theta)$$



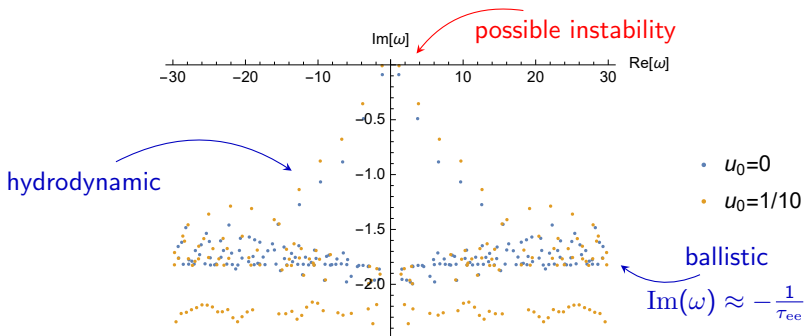
$$a_n = (-1)^n a_{-n}$$

“Dirty” (uniformly outgoing)
boundary conditions:



Eigenvalue spectrum ω (clean BCs)

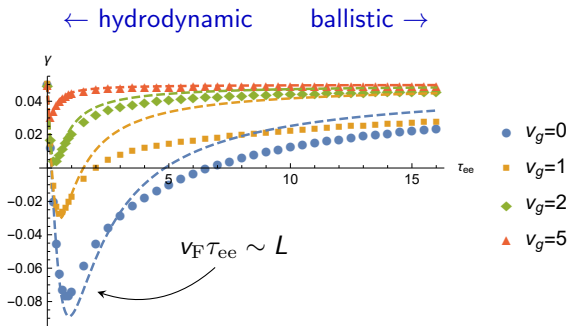
Dyakonov-Shur instability: $\gamma > 0$ with $\gamma = \max(\text{Im}(\omega))$



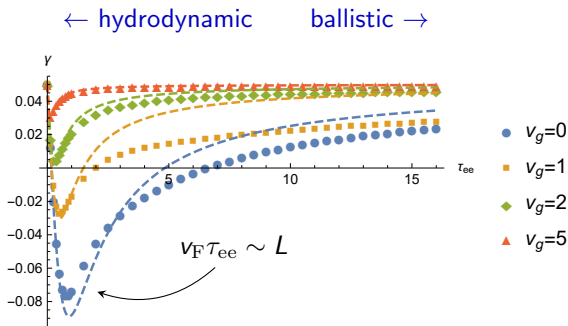
Parameters: $v_g = 0$, $\tau_{ee} = 1/2$, clean BCs

Mendl and Lucas, arXiv:1801.01501 (2018)

γ as a function of τ_{ee} (clean BCs)



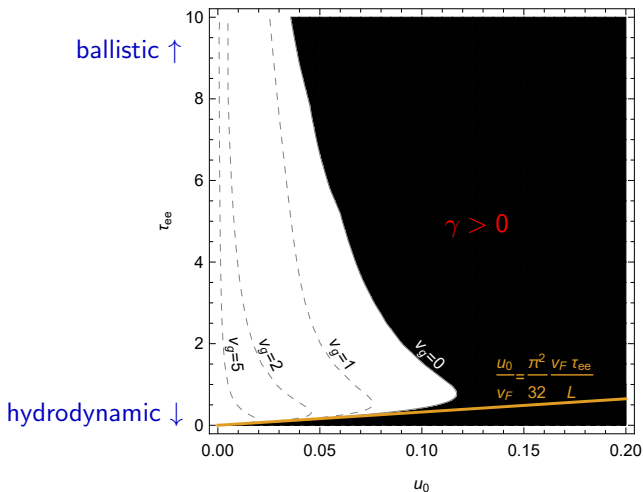
γ as a function of τ_{ee} (clean BCs)



Dashed: heuristic formula $\gamma \approx \frac{u_0}{L} - \frac{\pi^2 \nu}{8L^2} \frac{1}{1 + \left(\frac{\pi v_s}{2L} \tau_{ee}\right)^2}$

$$v_s = \sqrt{\frac{v_F^2}{2} + v_g^2} \quad \left| \begin{array}{l} \text{sound velocity} \\ \text{viscosity} \end{array} \right.$$
$$\nu = \frac{v_F^2 \tau_{ee}}{4}$$

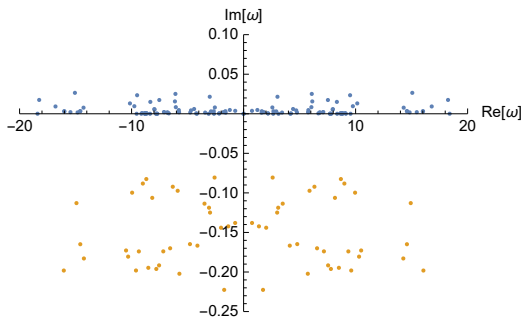
Instability ($\gamma > 0$) depending on τ_{ee} and u_0 (clean BCs)



Orange: $\gamma = 0$ contour for $\gamma = \frac{u_0}{L} - \frac{\pi^2 \nu}{8L^2}$

Eigenvalue spectrum in the ballistic limit, clean vs. dirty

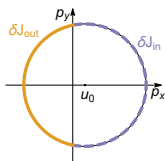
Ballistic limit $\tau_{ee} \rightarrow \infty$ (no collisions)



- clean BC
- noslip-dirty BC

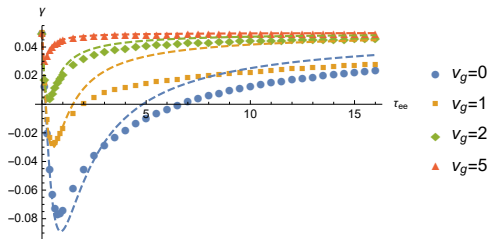
Parameters: $u_0 = 1/10$

Dirty boundary conditions:













Summary and conclusion

- Fate of the instability in the ballistic limit sensitive to boundary conditions
- Hydrodynamic regime shrinks substantially if $v_g \gg v_F \rightsquigarrow$ DS instability need not be hydrodynamic



- Absence of the DS instability in experiments could be used as a heuristic *upper and lower bound* on τ_{ee} and ν

References

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